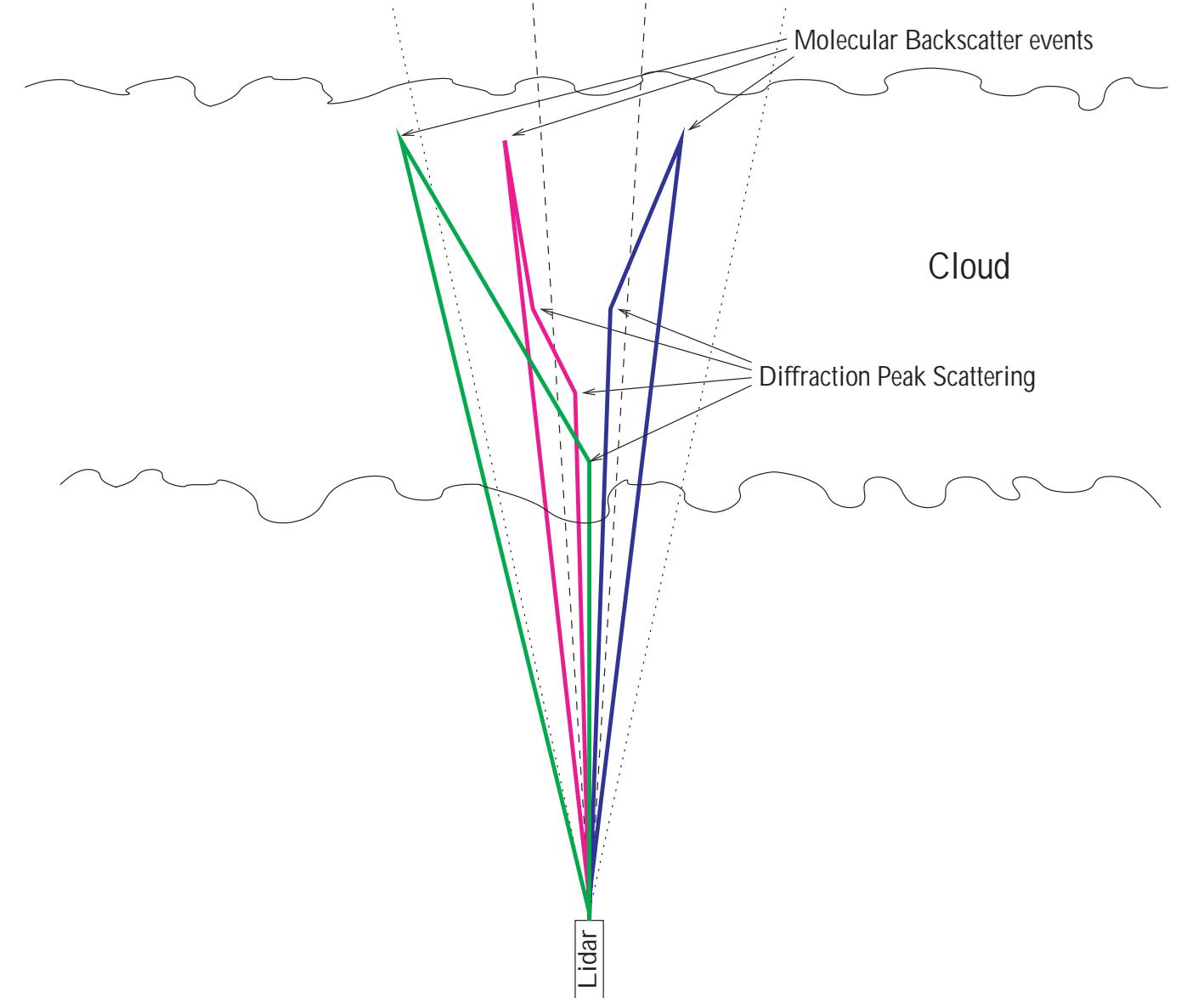
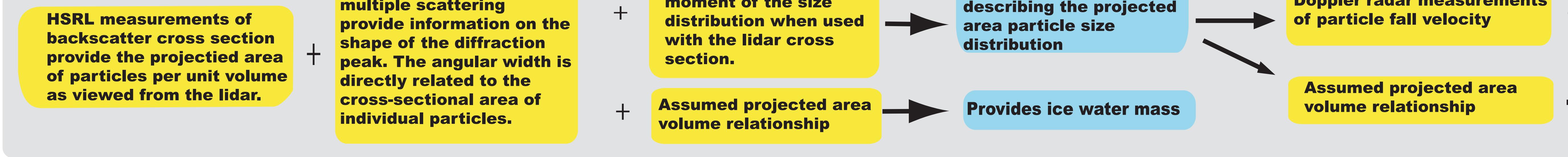


Lidar Multiple Scattering Determinations of Particle Size in Cirrus Clouds

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Basic Approach



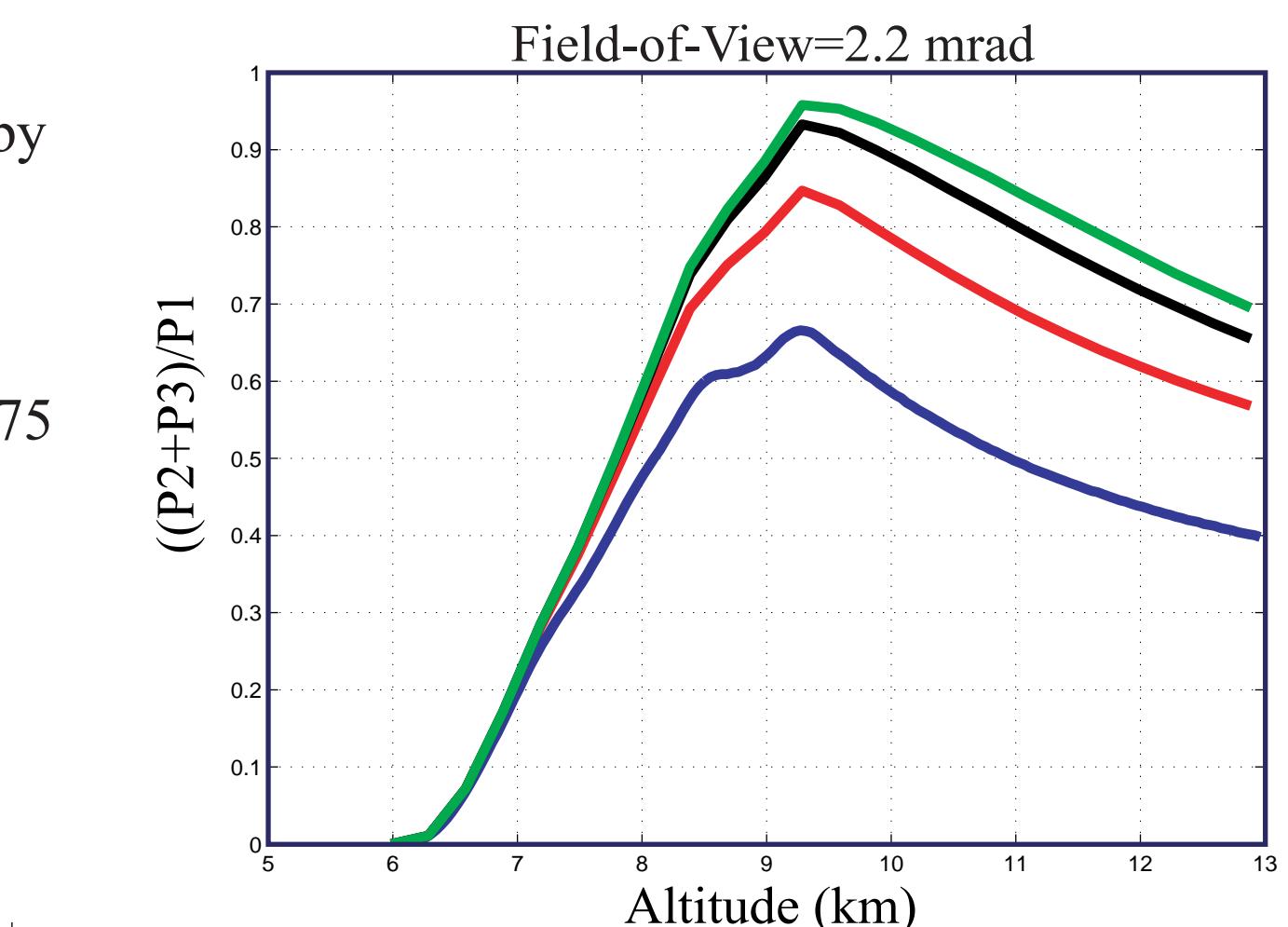
Exponential Distribution

$$\frac{dn}{dr} = a \exp(-br)$$

$$b = \frac{3}{r_{\text{eff}}}$$

Optical-radar radius = effective radius
 $<r> = 1.45 r_{\text{eff}}$

$$<r> = \sqrt{\frac{r^6}{<r^2>}}$$



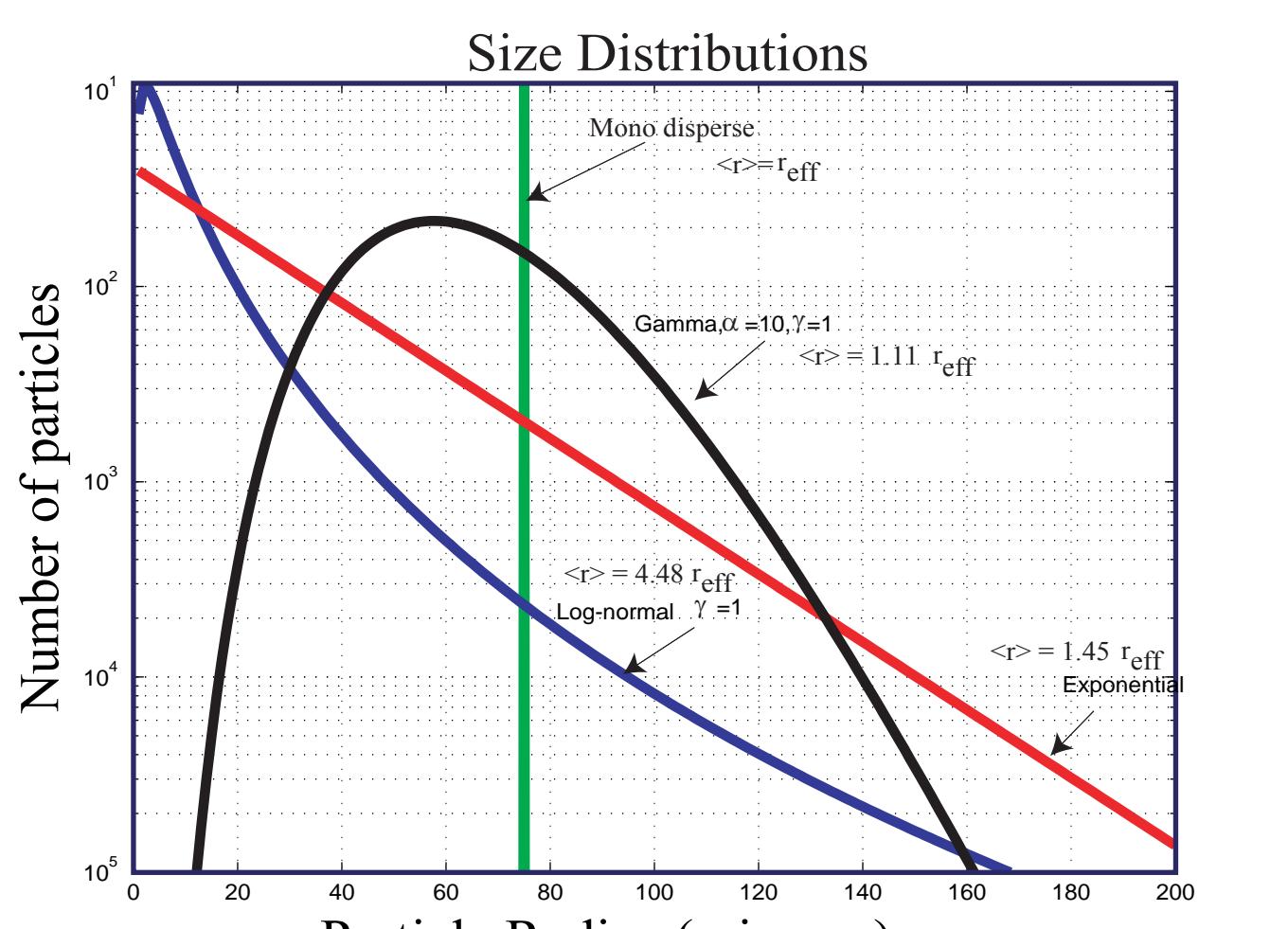
Changes in lidar multiple scattering caused by changing the shape of the size distribution. Multiple scattering was computed using the scattering cross section profile measured on 22-Feb-01 with the effective radius fixed at 75 microns. The sum of 2nd and 3rd order scattering divided by first order scattering is shown for each size distribution and each receiver field-of-view. Size distributions are designated by color.

Log-Normal Distribution

$$\frac{dn}{dr} = \frac{a}{\alpha r} \exp\left[-\frac{1}{2}\left(\frac{\ln(r/\alpha)}{\gamma}\right)^2\right]$$

$$\alpha = r_{\text{eff}} \exp(-2.5\gamma^2)$$

Optical-radar radius = effective radius
 $<r> = r_{\text{eff}} \exp(1.5\gamma^2)$

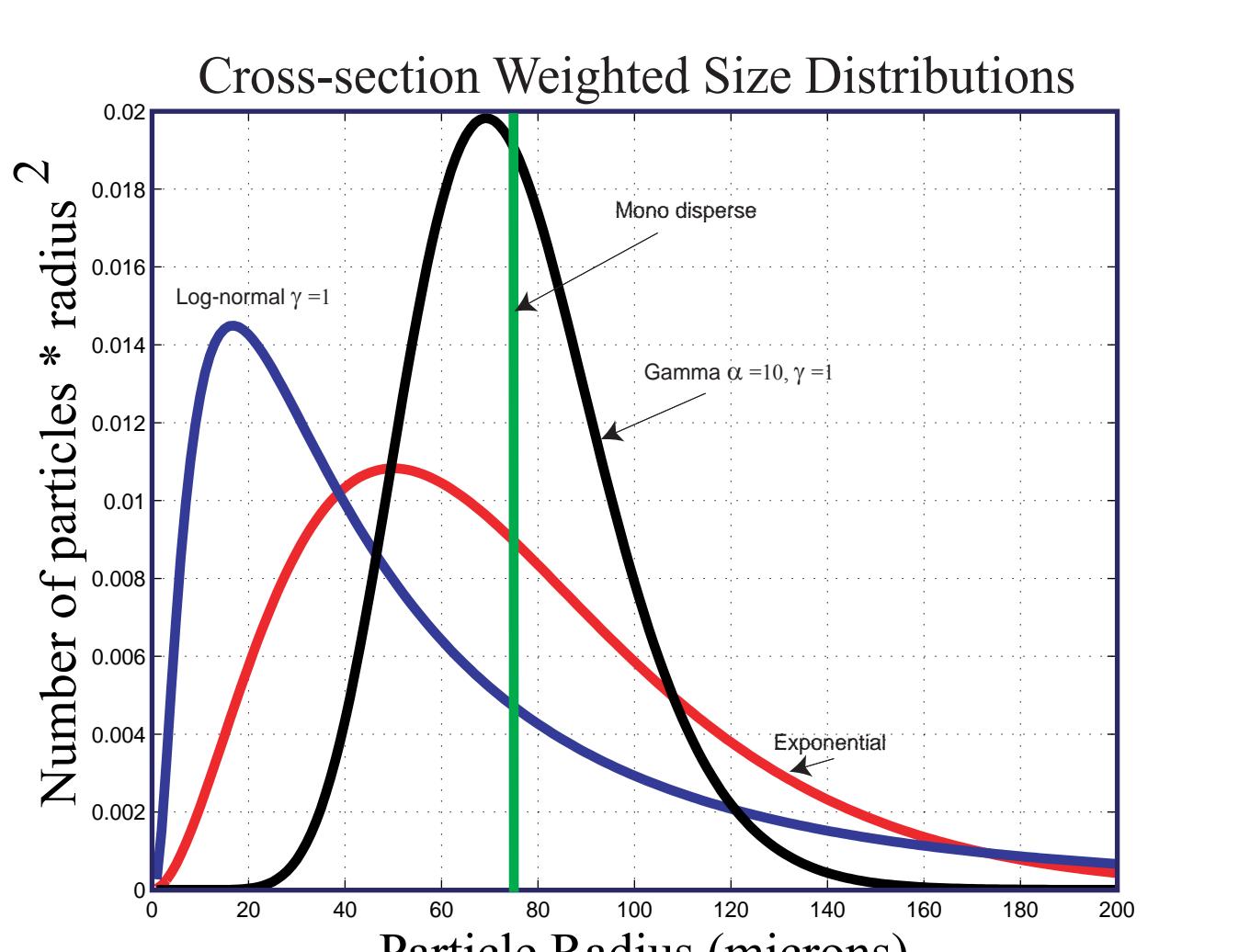


Gamma Distribution

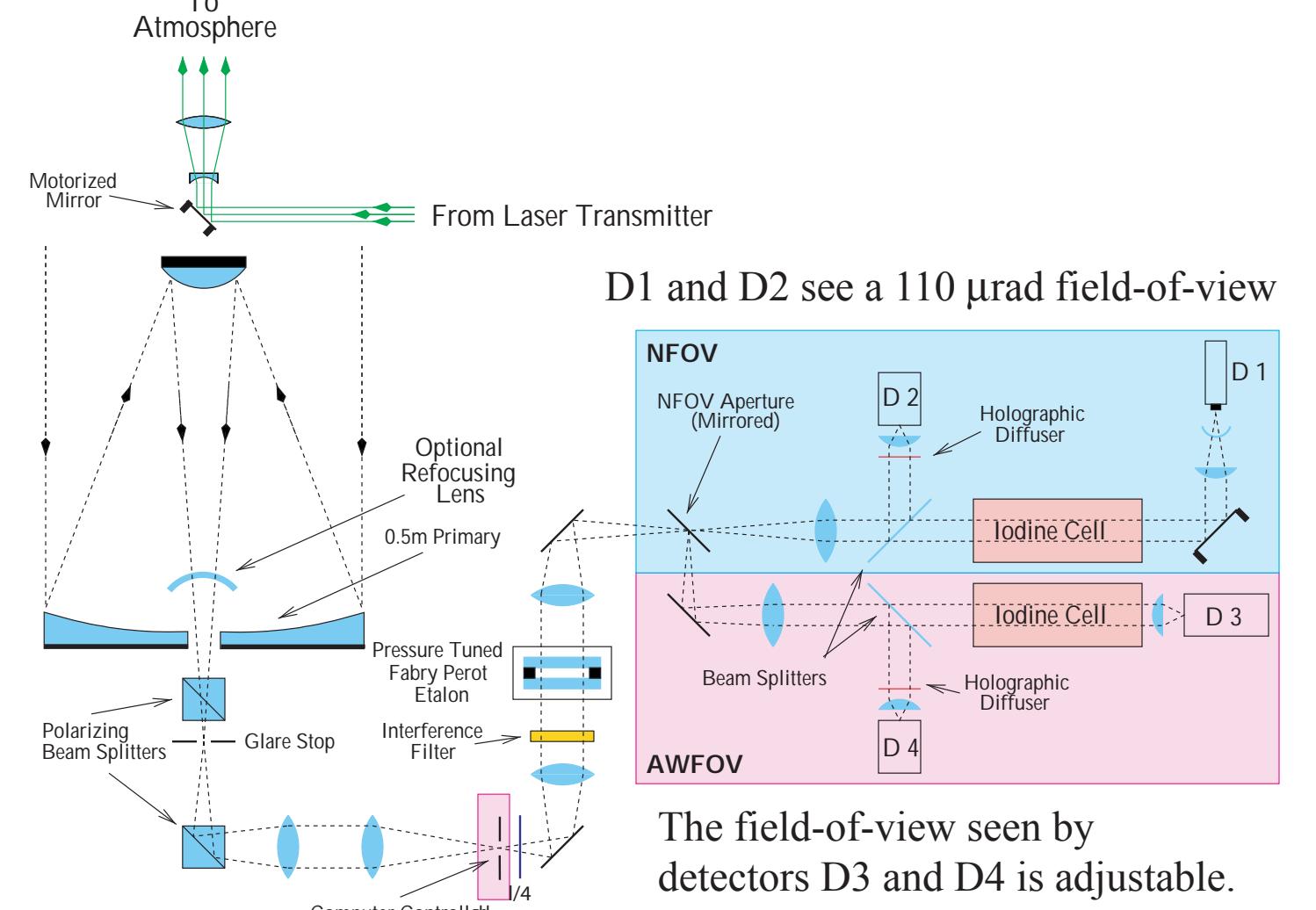
$$\frac{dn}{dr} = a r^\alpha \exp(-b r^\gamma)$$

$$b = \frac{1}{r_{\text{eff}}} \frac{\Gamma(\alpha+4)}{\Gamma(\alpha+3)}$$

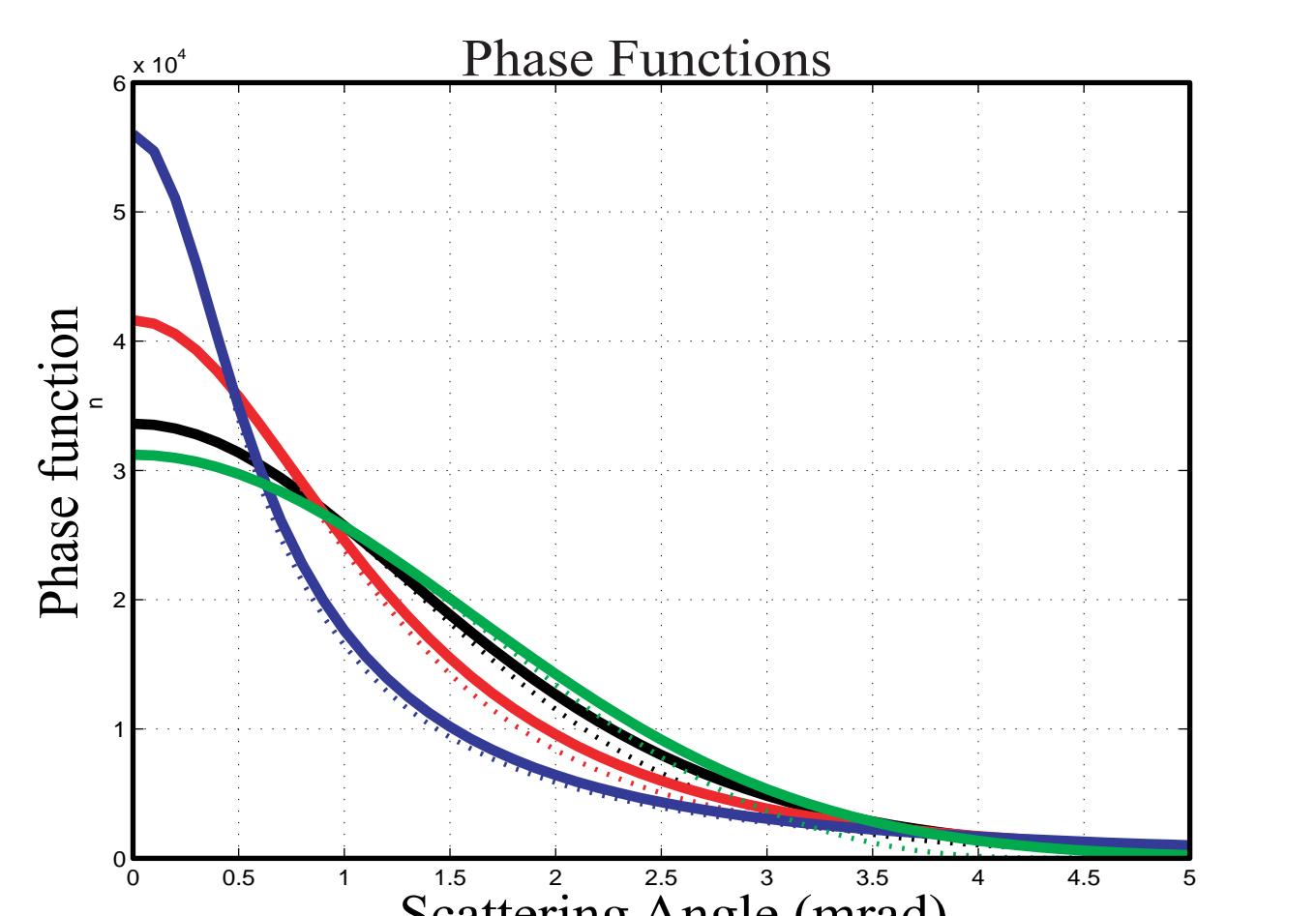
Optical-radar radius = effective radius
 $<r> = r_{\text{eff}} \frac{\Gamma(\alpha+3)}{\Gamma(\alpha+4)}$



HSRL multiple field of view measurements and a multiple scattering model provide information on forward diffraction peak of the scattering phase function. This is used to derive particle size distribution parameters. This distribution describes the dimensions of particles projected on a plane perpendicular to the lidar beam. The HSRL is able to isolate photons which have undergone one or more small angle forward scatterings coupled with one molecular backscatter event.



HSRL Receiver Schematic. Geiger-mode APD (D1), and wide-field-of-view channels have been made operational under this grant. The APD has provided a factor of 10 improvement in the sensitivity of this channel.



A comparison of the phase functions for the different size distributions using the Gaussian approximation (Solid lines), and using diffraction theory (dashed lines).

We have derived multiple scatter lidar equations describing the ratio of nth order multiple scattering to first order scattering as a function of range (R) for several mathematical models of the particle size distribution. These equations assume :

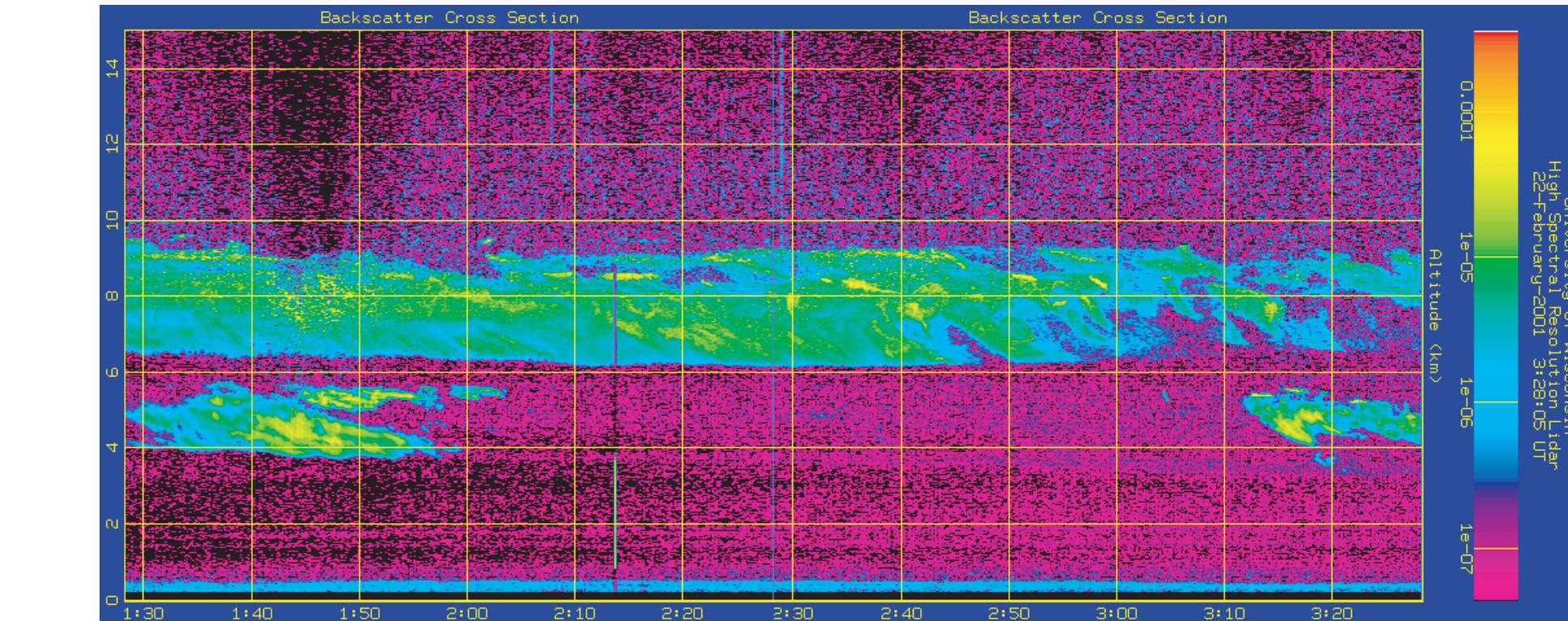
- 1) Particle sizes described by log-normal, exponential, gamma or mono disperse distributions.
- 2) Particles that are large compared to the lidar wavelength, λ .
- 3) A Gaussian distribution of energy in the transmitted beam with an angular width = $2\rho_1$.
- 4) A backscatter phase function $P(\pi, R)$ with an average value of $P_{\text{nt}}(R)$ for angles near π .
- 5) A cloud optical depth = $\tau(R)$, scattering cross section = $\beta(R)$

For the log-normal distribution:

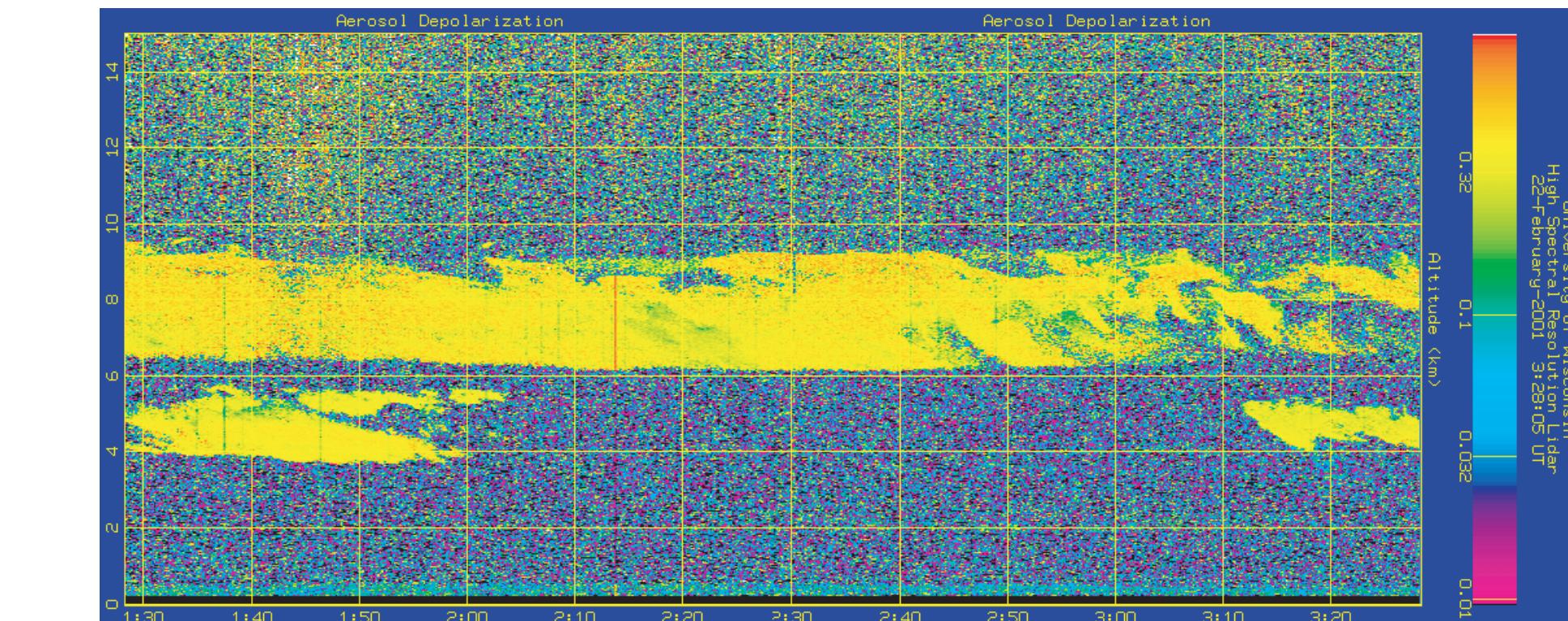
$$\frac{P_n(R)}{P_1(R)} = \frac{P_{\text{nt}}(R)}{P(\pi, R)} \left[1 - \exp\left(-\frac{R^2}{\rho_1^2}\right) \right]^{-1} \left[\frac{\tau^{n-1}}{(n-1)!} \int_{z_1}^R \frac{\beta(z_1)}{\sqrt{\pi}} \int_{z_2}^R \exp(-u^2) \cdots \int_{z_{n-1}}^R \frac{\beta(z_{n-1})}{\sqrt{\pi}} \int_{z_n}^R \exp(-u^2) \exp\left(-\frac{-\pi^2 \rho_1^2 R^2 / \lambda^2}{(R-z_1)^2 \left(\exp(-\sqrt{2}\gamma(z_1)u_1 - 2\gamma(z_1)^2) \right) + \cdots (R-z_{n-1})^2 \left(\exp(-\sqrt{2}\gamma(z_{n-1})u_{n-1} - 2\gamma(z_{n-1})^2) \right) + \pi^2 \rho_1^2 R^2 / \lambda^2} \right) du'_1 dz_{n-1} du'_{n-2} dz_{n-2} \cdots du'_1 dz_1 \right]$$

For the gamma distribution:

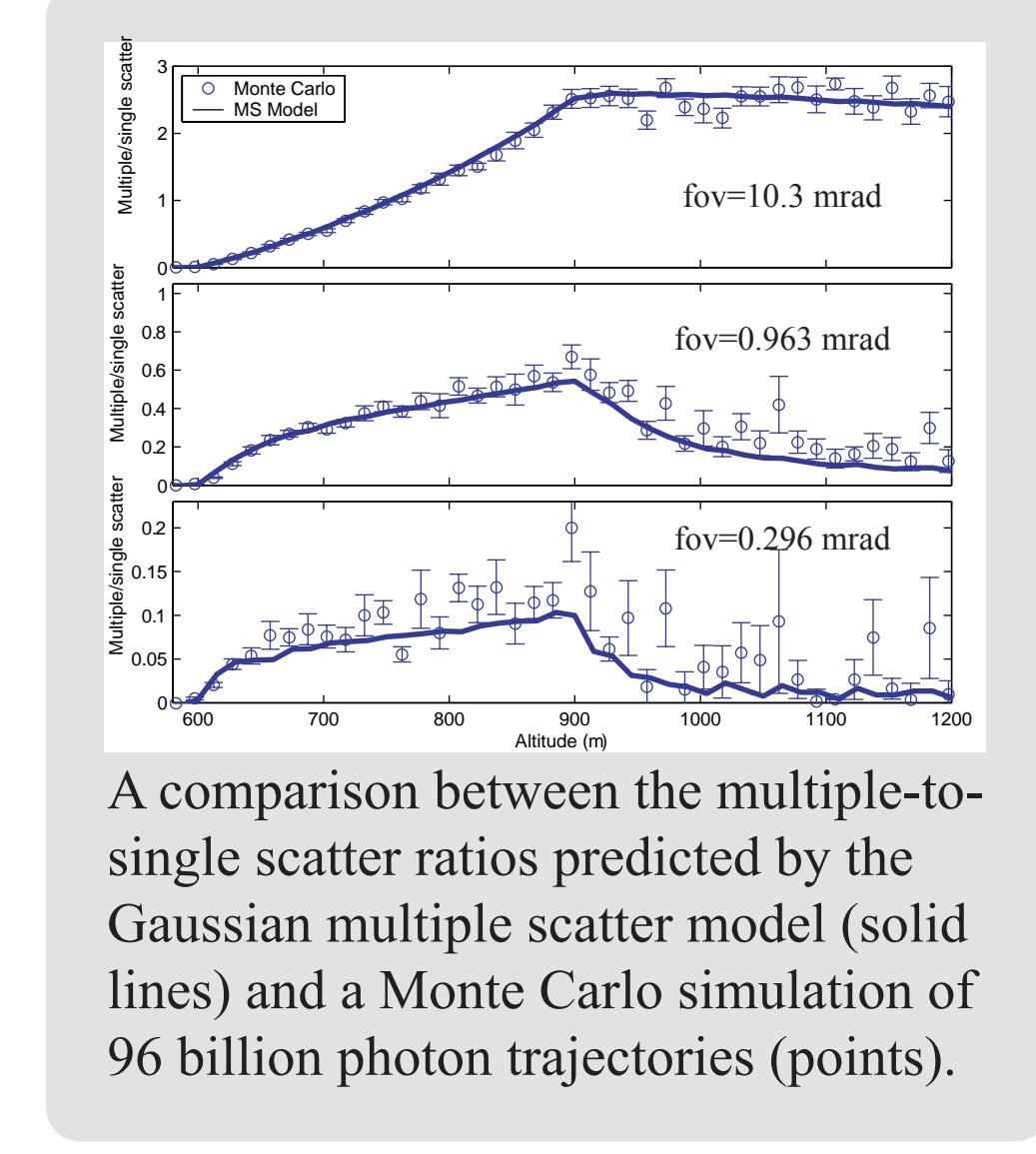
$$\frac{P_n(R)}{P_1(R)} = \frac{P(\pi, R)}{P(\pi, R)} \left[1 - \exp\left(-\frac{R^2}{\rho_1^2}\right) \right]^{-1} \left[\frac{\tau^{n-1}}{(n-1)!} \int_{z_1}^R \frac{\beta(z_1)}{\Gamma(\alpha z_1 + 3)} \int_{z_2}^R u(z_1)^{\alpha z_1 + 2} \exp(-u(z_1)^{\alpha z_1}) \int_{z_3}^R \beta(z_{n-1}) \frac{\gamma(z_{n-1})}{\Gamma(\alpha z_{n-1} + 3)} \int_{z_n}^R u(z_{n-1})^{\alpha z_{n-1} + 2} \exp(-u(z_{n-1})^{\alpha z_{n-1}}) \exp\left(-\frac{\pi^2 \rho_1^2 R^2 / \lambda^2}{(R-z_1)^2 \left(\frac{1}{u(z_1) r_{\text{eff}}(z_1)} \frac{1}{\Gamma(\alpha z_1 + 3)} \right)^2 + \cdots (R-z_{n-1})^2 \left(\frac{1}{u(z_{n-1}) r_{\text{eff}}(z_{n-1})} \frac{1}{\Gamma(\alpha z_{n-1} + 3)} \right)^2 + \pi^2 \lambda^{-2} \rho_1^2 R^2} \right) du'_1 dz_{n-1} du'_{n-2} dz_{n-2} \cdots du'_1 dz_1 \right]$$



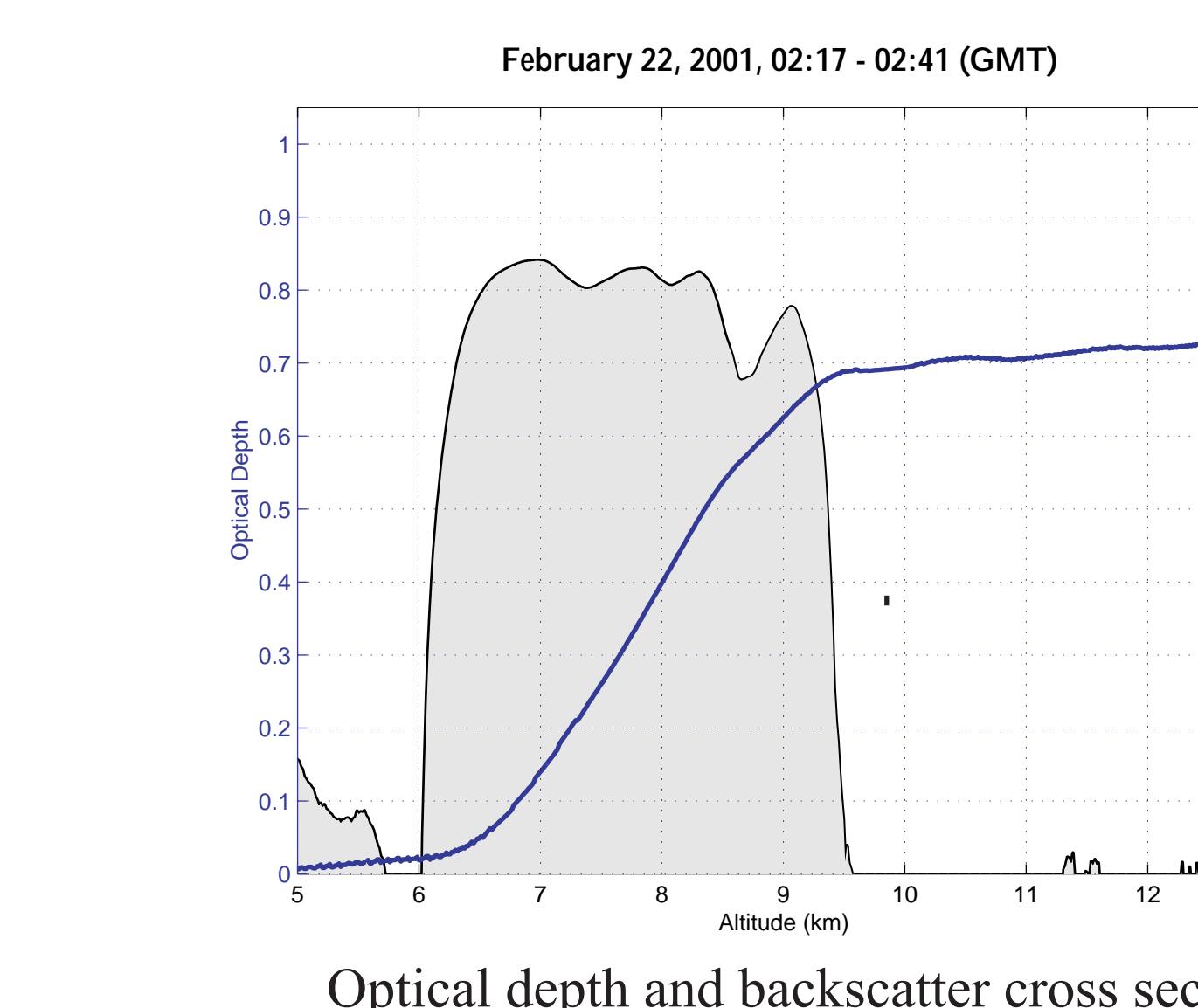
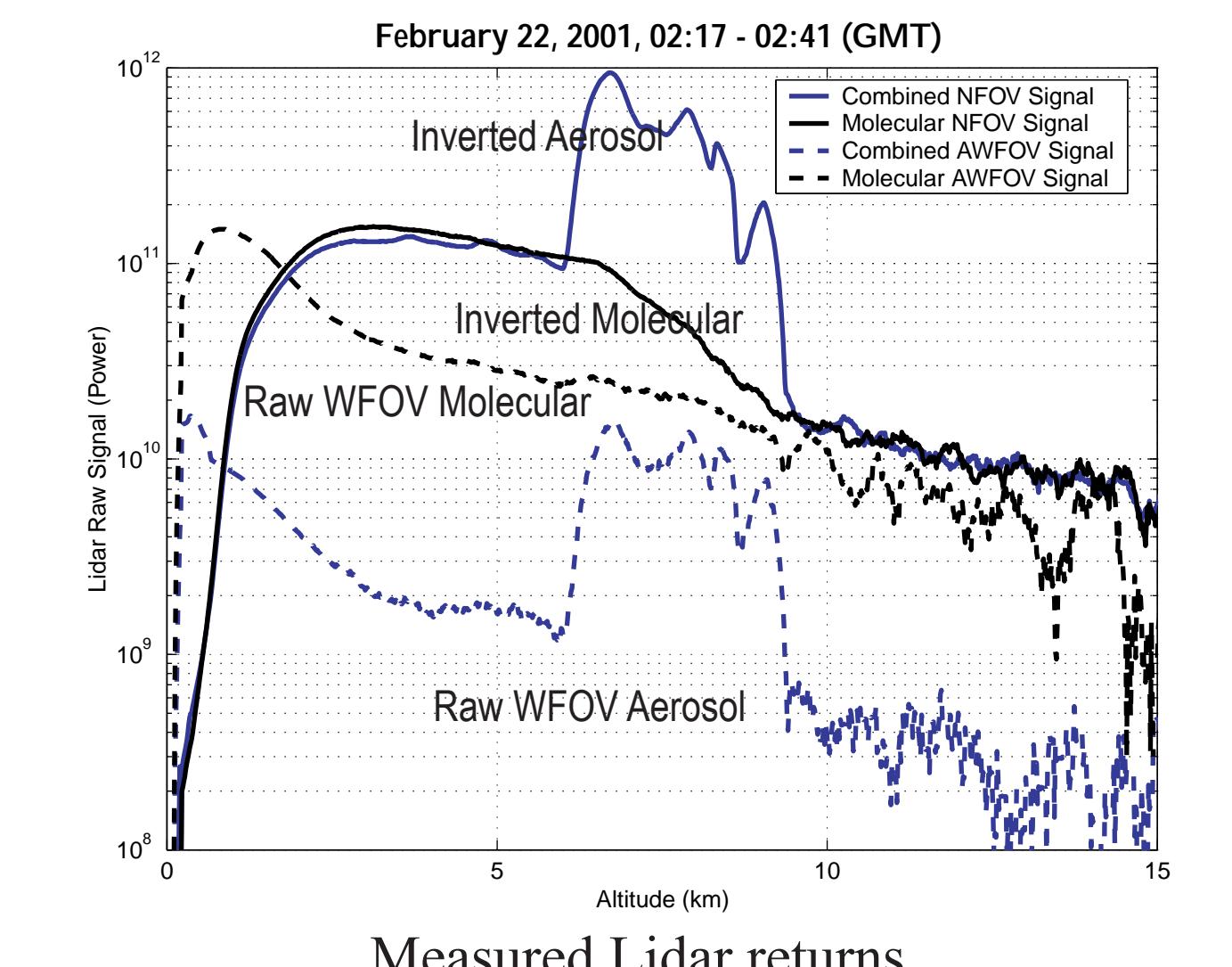
Cirrus backscatter cross section 22-Feb-01



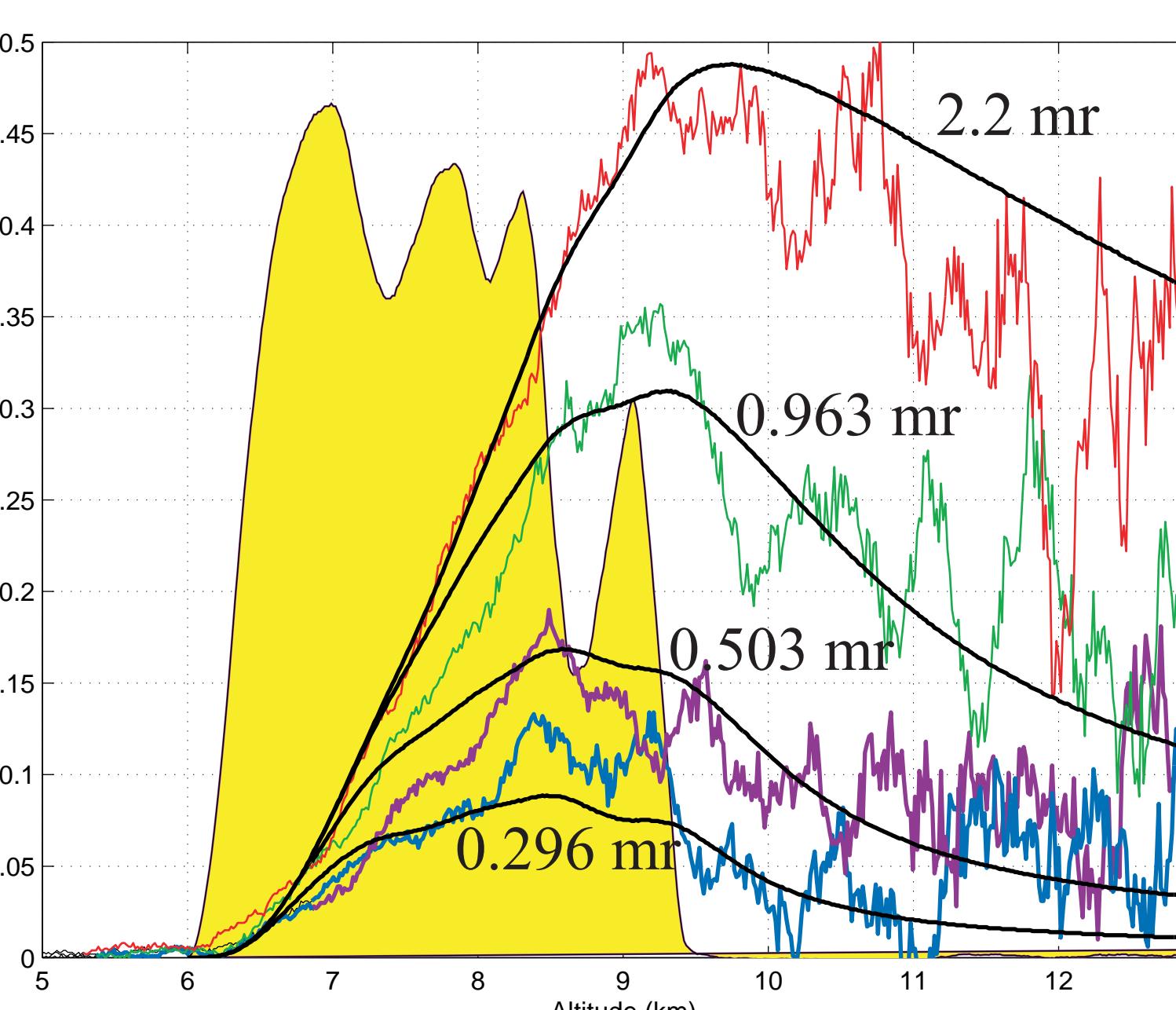
Cirrus cloud depolarization, 22-Feb-01



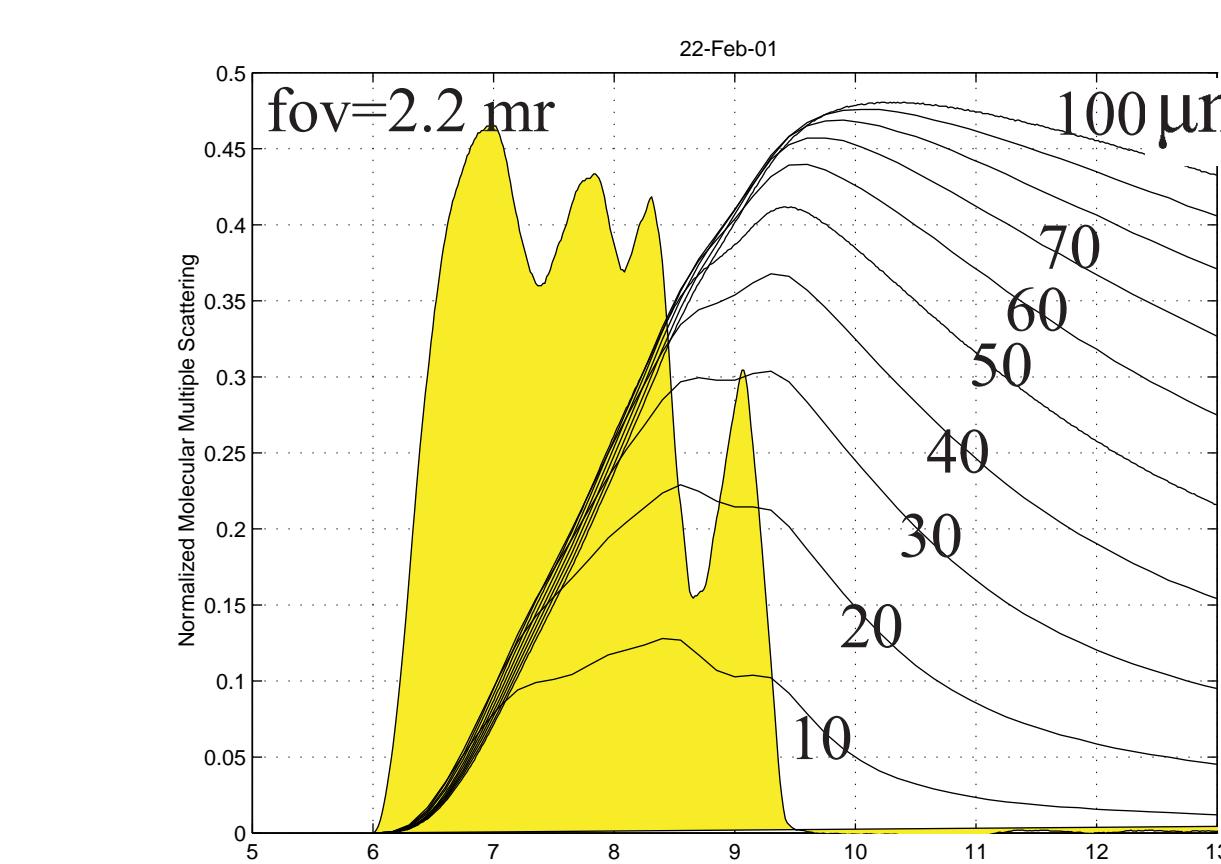
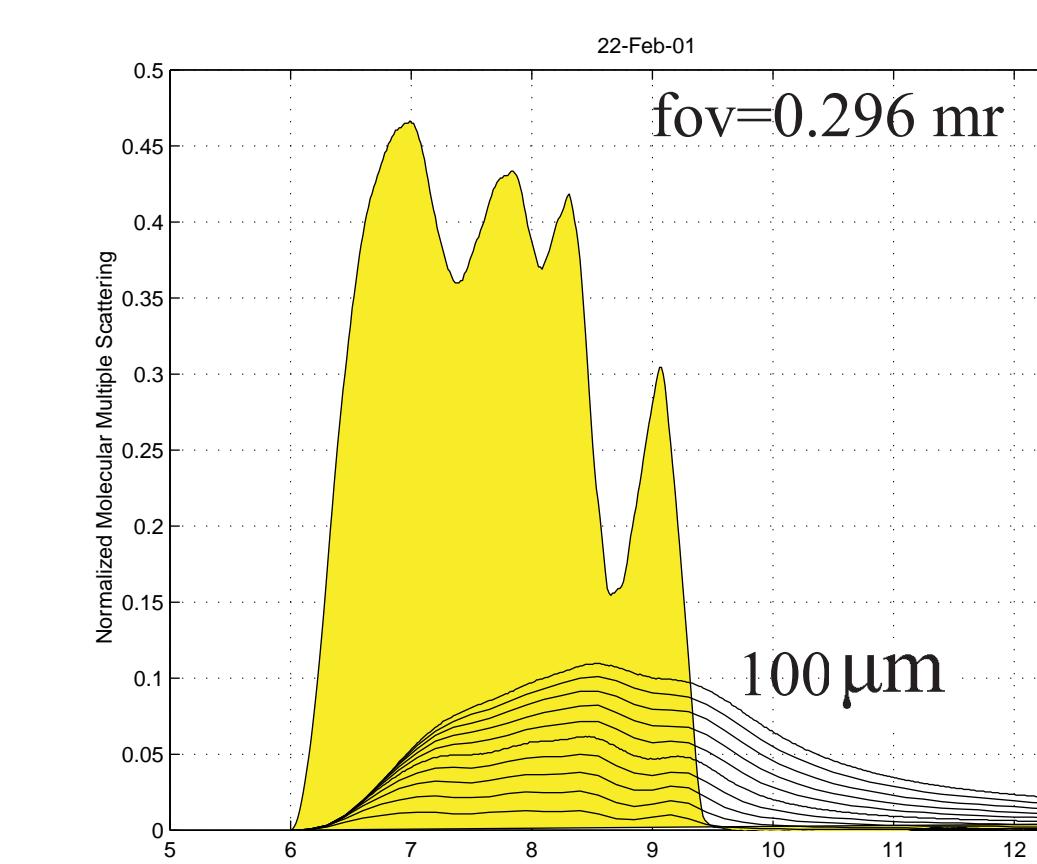
A comparison between the multiple-to-single scatter ratios predicted by the Gaussian multiple scatter model (solid lines) and a Monte Carlo simulation of 96 billion photon trajectories (points).



Optical depth and backscatter cross section.



A comparison of the measured multiply scattered molecular lidar return (colored lines) and the model predictions (solid black lines) using $r_{\text{eff}} = 75$ microns and $\gamma = 0$. The backscatter cross section is shown in yellow.



The normalized molecular wide field of view lidar return computed as function of receiver field of view and r_{eff} with $\gamma = 0$. The backscatter cross section profile of the cloud is shown in yellow (relative units).

$$\text{WFOV signal-}k^*(110 \text{ urad FOV molecular signal}) \\ (110 \text{ urad FOV molecular signal})$$

$$\text{Where: } k = \frac{\text{Clear air WFOV Signal}}{\text{Clear air 110 urad FOV signal}}$$

and where WFOV is the molecular wide-field of view channel.